

Nonthermal Distribution of Electro-negative Light-ions in Heavy-ion-acoustic Solitary and Shock Structures

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Abstract An attempt has been taken to study the nonlinear propagation of electro-positive heavy ion-acoustic (EPHIA) solitary and shock waves in multi-component plasmas containing mobile electro-positive heavy ions (EPHIs), nonthermal distributed electro-negative light ions (ENLIs) and Maxwellian electrons. The effects of temperature, mobile EPHIs, ENLIs and its nonthermality on solitary and shock waves are investigated in present work. To do so, we have derived the standard Korteweg-de Vries (K-dV) and Burger equations and their solutions. We have studied the effects of the parametric regimes, which affect the solitary and shock waves, numerically as well as graphically.

Keywords: Electro-positive heavy ion, solitary and shock waves, nonthermal distribution.

1 Introduction

Ion-acoustic (IA) solitary waves (SWs) are one of the important non-linear phenomena in plasma systems. Ions may be present in most of the plasmas and the presence of ions leads to existence of several low-frequency waves. IA SWs in unmagnetized plasma have been studied by a number of authors both experimentally and theoretically. Ions are considered as electro-positive, in most of the earlier works, but Conney *et al.* [1] found the existence of electro-positive (EP) and electro-negative (EN) ions plasma in their experimental investigation. The latter may also support the formation of IA shocks when the ratio of the negative ion to positive ion number density exceeds about 0.9 [2]. The EP and EN ions in plasmas are found in natural and technological environments like the D-region of the Earth's ionosphere, the Earth's mesosphere, the solar photosphere, and the microelectronics plasma processing reactors [3]. In the recent few years, the study of Korteweg-de Vries (K-dV) and Burger ion-acoustic solitons and shocks in a multi-species plasma consisting of electro-positive ions (EPIs), electro-negative ions (ENIs), and electrons become the field of current research investigation. Mishra *et al.* [4] studied the obliquely propagating ion-acoustic solitary waves in a multi-component magnetized plasma consisting of warm adiabatic positive and negative ion species and hot isothermal electrons. Haider *et al.* have studied the nonlinear propagation of IA solitary waves with EPIs and ENIs for Maxwellian [6,5], trapped [7,8], and nonthermal [9,10] distributed electrons. Nakamura and Tsukabayashi [11] have studied experimentally the propagation of ion-acoustic solitons in a plasma with negative ions.

The effects of heavy ions on radio wave propagation in the ionosphere have been studied by Hines [12] considering a multi-ion plasma around sixty years ago. Recently Hossen *et al.* [13] showed that heavy ions significantly modify the basic features of heavy-ion-acoustic (HIA) waves in degenerate plasmas. Ema *et al.* [14] have studied the dust-ion-acoustic (DIA) shock structures in multi-ion unmagnetized plasmas containing negatively charged heavy ions, positively charged light ions following Maxwellian distribution, and nonextensive electrons. Shah *et al.* [15] have theoretically investigated properties of HIA solitary and shock waves in unmagnetized plasmas with Maxwellian light ions. All of the above investigations, light ions are considered as Maxwellian. Very recently Rahaman and Haider [16] considered that ENLIs as trapped. But the ENLIs might very cold and following the nonthermal distribution. Therefore a necessity is felt to study the effect of nonthermal ENLIs in electro-positive HIA (EPHIA) solitary and shock waves in multi-ion plasmas. In the present work, the propagation of EPHIA solitary structures have been studied in unmagnetized plasma consisting of EPHIs, ENLIs following nonthermal distribution and Maxwellian

distributed electrons. The reductive perturbation method [17] has been employed to derive the K-dV [18] and Burgers [19] equations.

The rest of the manuscript is organized as follows. The plasma model for a small but finite amplitude EPHIA solitary and shock waves in an unmagnetized collisionless multi-component plasma with nonthermal ENLIs is presented in section 2. The well known K-dV equation is derived with its solution to study the solitary waves in section 3. To study the shock waves, Burgers equation is derived as well as its solution in section 4. The numerical study of this work is presented in section 5. Finally, a summary of our results is discussed in section 6.

2 Plasma Model

We have considered a collisionless, unmagnetized multi-ion plasma system where -

- i) EPHIs are mobile,
- ii) ENLIs are nonthermal,
- iii) Electrons follow the Maxwellian distribution.

The one dimensional dynamics (equation of continuity, momentum equation and Poisson's equation) of the EPHIA waves in the normalized form are governed by

$$\frac{\partial n_p}{\partial t} + \frac{\partial}{\partial x}(n_p u_p) = 0 \quad (1)$$

$$\frac{\partial u_p}{\partial t} + u_p \frac{\partial u_p}{\partial x} = -\frac{\partial \phi}{\partial x} + \eta \frac{\partial^2 u_p}{\partial x^2} \quad (2)$$

$$\frac{\partial^2 \phi}{\partial x^2} = \mu_e n_e + \mu_n n_n - n_p \quad (3)$$

where, $n_p(n_n)$ is the EPHIs (ENLIs) number density normalized by its equilibrium value $n_{p0}(n_{n0})$, u_p is the positive ion fluid speed normalized by $C_p = (K_B T_e / m_p)^{1/2}$ with K_B is the Boltzmann constant, T_e is the temperature of electrons and m_p is the rest mass of EPHIs. ϕ is the EPHIA wave potential normalized by $K_B T_e / e$ with e being the magnitude of the charge of the electron. The time variable (t) is normalized by $\omega_p^{-1} = (4\pi n_{p0} e^2 / m_p)^{1/2}$ with c being the speed of light. The space variables are normalized by Debye radius $\lambda_D = (K_B T_e) / (4\pi n_{p0} e^2)$. We have considered the viscous term, i.e. coefficient of viscosity (η) which equals to zero at the time of studying solitary waves.

Maxwillan electron distribution can be express as

$$n_e = e^\phi$$

According to Cairns et al. [20] the nonthermal distribution of ENLIs with a population of fast particles (γ) can be represented as

$$n_n = [1 - \sigma_p \alpha \phi + \alpha (\sigma_p \phi)^2] e^{(\sigma_p \phi)}$$

where $\alpha = 4\gamma / (1 + 3\gamma)$ with γ represents a parameter determining the fast particles present in this plasma model and $\sigma_p (= \frac{T_e}{T_n})$ is the temperature ratio of electron (T_e) to ENLIs (T_n).

Now using equilibrium charge neutrality condition $n_{e0} + n_{n0} = n_{p0}$, one can write $\mu_e = 1 - \mu_n$ where, $\mu_e = \frac{n_{e0}}{n_{p0}}$ and $\mu_n = \frac{n_{n0}}{n_{p0}}$.

3 Solitary Wave

To derive the Korteweg-de Vries (K-dV) equation from the normalized equations (1)-(3) by employing the reductive perturbation technique, we have introduced stretched co-ordinates as [7,21],

$$\xi = \epsilon^{1/2} (x - v_p t)$$

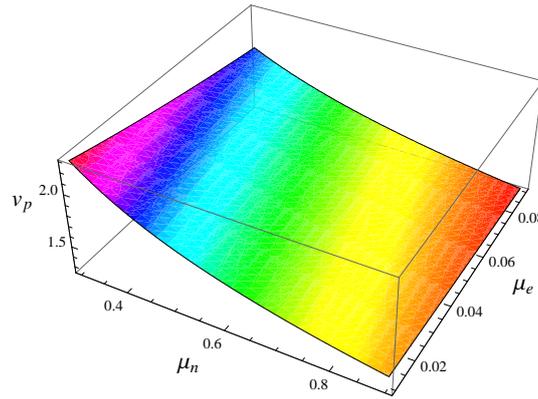


Figure 1. Variation of the phase speed with μ_n and μ_e for $u_0 = 0.1$, $\alpha = 0.7$, and $\sigma_p = 2$.

$$\tau = \epsilon^{\frac{3}{2}} t$$

where ϵ is the smallness parameter measuring the weakness of the amplitude or dispersion and v_p is the phase speed of the system normalized by C_p .

The perturbed quantities can be expanded about their equilibrium values in powers of ϵ as

$$\left. \begin{aligned} n_p &= 1 + \epsilon n_p^{(1)} + \epsilon^2 n_p^{(2)} + \dots \\ u_p &= 0 + \epsilon u_p^{(1)} + \epsilon^2 u_p^{(2)} + \dots \\ \phi &= 0 + \epsilon \phi^{(1)} + \epsilon^2 \phi^{(2)} + \dots \end{aligned} \right\} \quad (4)$$

Using the stretched coordinates and Eq.(4) in Eqs. (1)-(3), one can obtain the following linear dispersion relation by taking the lowest order of Eqs. (1)-(3).

$$v_p = \frac{1}{\sqrt{\mu_e + \mu_n \sigma_p (1 - \alpha)}} \quad (5)$$

To get next higher order equations, we have substituted the required parameter values from the above equations into Eqs.(1)-(3).

$$\frac{\partial n_p^{(1)}}{\partial \tau} - v_p \frac{\partial n_p^{(2)}}{\partial \xi} + \frac{\partial}{\partial \xi} (n_p^{(1)} u_p^{(1)}) + \frac{\partial u_p^{(2)}}{\partial \xi} = 0 \quad (6)$$

$$\frac{\partial u_p^{(1)}}{\partial \tau} - v_p \frac{\partial u_p^{(2)}}{\partial \xi} + u_p^{(1)} \frac{\partial u_p^{(1)}}{\partial \xi} = - \frac{\partial \phi^{(2)}}{\partial \xi} \quad (7)$$

$$\frac{\partial^2 \phi^{(1)}}{\partial \xi^2} = \frac{1}{v_p^2} \phi^{(2)} - n_p^{(2)} + \frac{1}{2} [\phi^{(1)}]^2 (\mu_e + \sigma_p^2 \mu_n) \quad (8)$$

Now using Eqs.(6) -(8), KdV equation can be readily obtain as

$$\frac{\partial \phi^{(1)}}{\partial \tau} + A \phi^{(1)} \frac{\partial \phi^{(1)}}{\partial \xi} + B \frac{\partial^3 \phi^{(1)}}{\partial \xi^3} = 0 \quad (9)$$

where nonlinear and dissipation coefficients respectively are

$$A = \frac{3}{2v_p} - \frac{(\mu_e + \sigma_p^2 \mu_n) v_p^3}{2} \quad (10)$$

$$B = \frac{v_p^3}{2} \quad (11)$$

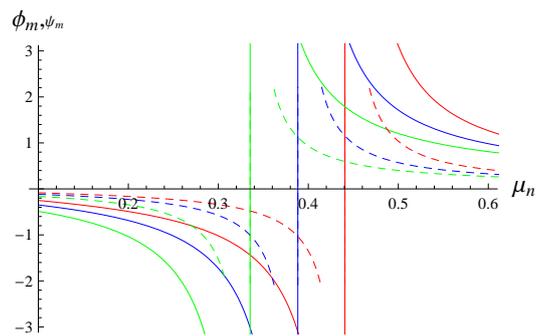


Figure 2. Variation of the amplitude of solitary (solid curve, ϕ_m) and shock (dashed curve, ψ_m) waves with respect to μ_n for $\sigma_p = 2$, $\alpha = 0.7$, and $u_0 = 0.1$. Here red curve represents $\mu_e = 0.01$, blue curve is for $\mu_e = 0.05$, and green curve is for $\mu_e = 0.09$

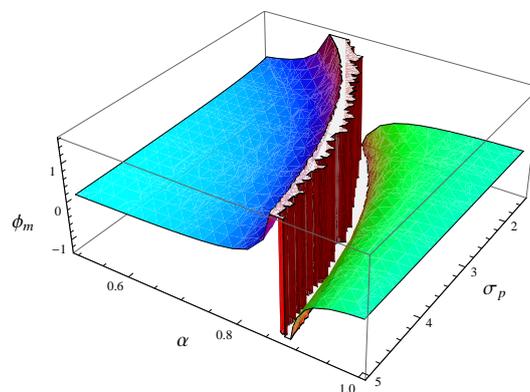


Figure 3. Variation of the amplitude of solitary wave (ϕ_m) with respect to α and σ_p for $\mu_e = 0.05$, $\mu_n = 0.5$, and $u_0 = 0.1$.

The stationary solitary wave solution [5,10,22,23] of the K-dV equation is

$$\phi = \phi_m \operatorname{sech}^2 \left[\frac{(\xi - u_0\tau)}{\delta_1} \right] \tag{12}$$

where u_0 is the constant SW velocity. The amplitude (ϕ_m) and the width (Δ) of the solitary wave equation are

$$\phi_m = \frac{3u_0}{A} \tag{13}$$

$$\delta_1 = \sqrt{\frac{4B}{u_0}} \tag{14}$$

4 Shock Wave

To obtain Burgers equation using reductive perturbation method we have introduced the following stretched co-ordinates [22,23]

$$\xi = \epsilon(x - v_p t)$$

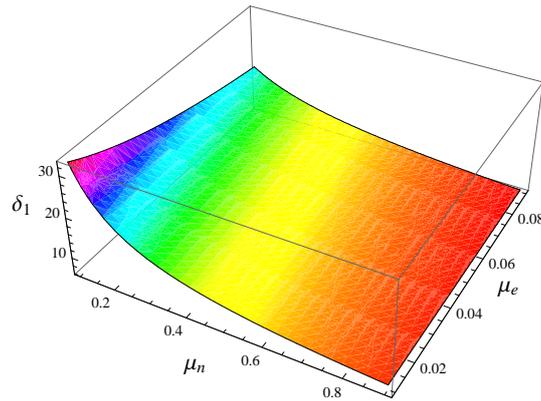


Figure 4. Variation of the width of solitary wave (δ_1) with respect to μ_n and μ_e for $\sigma_p = 2$, $\alpha = 0.7$, and $u_0 = 0.1$.

$$\tau = \epsilon^2 t$$

Now, using the same procedure, done above (in Sec. 3), with the new stretched co-ordinates (for shock wave) we can obtain the linear dispersion relation similar as solitary waves, as in (5).

By following the same algebraic calculation as in Sec. 3, one can be readily obtain Burgers equation as

$$\frac{\partial \phi^{(1)}}{\partial \tau} + A \phi^{(1)} \frac{\partial \phi^{(1)}}{\partial \xi} - C \frac{\partial^2 \phi^{(1)}}{\partial \xi^2} = 0 \tag{15}$$

where, nonlinear coefficient A is the same as Eq.(10) and

$$C = \frac{\eta}{2} \tag{16}$$

The stationary solution of the Burgers equation (15) can express as [22,23]

$$\psi = \psi_m [1 - \tanh \frac{(\xi - u_0 \tau)}{\delta_2}] \tag{17}$$

where the amplitude and width are given by

$$\psi_m = \frac{1}{3} \phi_m \tag{18}$$

$$\delta_2 = \frac{2C}{u_0} \tag{19}$$

5 Numerical Study

We have numerically studied the effect of non-thermal ENLIs on EPHIA solitary and shock waves in multi-ion plasmas with Maxwellian electrons. It is obvious from Eqs.(13), (14), (18), and (19) that u_0 is a linear function of amplitude of both solitary and shock waves and inverse function of width which means that the profile of the faster solitary and shock structure will be taller and narrower than slower one. Figure 1 shows the variation of phase speed (v_p) with ENLIs concentration (μ_n) and electron concentration (μ_e) keeping the values $u_0 = 0.1$, $\alpha = 0.7$, and $\sigma_p = 2$. It is found that phase velocity decreases with increasing μ_n and μ_e . This is to note here that, at the critical value ($\gamma = 1$) the phase velocity v_p reduce to $1/\sqrt{\mu_e}$, which means that at critical condition, the phase velocity is independent of all the parameters except μ_e .

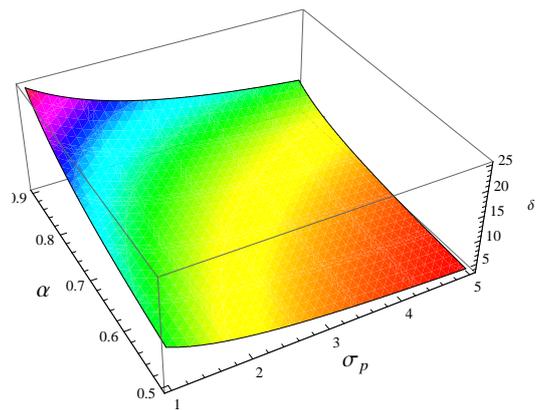


Figure 5. Variation of the width of solitary wave (δ_2) with respect to σ_p and α for $\mu_e = 0.05$, $\mu_n = 0.5$, and $u_0 = 0.1$.

Depending on the value of A the solitary and shock waves might be associated with positive or negative potentials. Equation (10) indicates that A is a function of σ_p , μ_n , μ_e and α . Therefore, these parameters are responsible for the solitary waves associated with positive and negative potentials. This is also shown in figs. 2 and 3. Figure 2 shows the variation of the amplitude of solitary (solid curve, ϕ_m) and shock (dashed curve, ψ_m) waves with μ_n keeping the values $u_0 = 0.1$, $\sigma_p = 2$, and $\alpha = 0.7$. The red curve corresponds to $\mu_e = 0.01$, blue curve to $\mu_e = 0.05$ and green curve to $\mu_e = 0.09$. It is observed from the figure that there is a critical value of μ_n considering other parameters constant below (above) which the solitary and shock waves are associated with negative (positive) potentials. Similar types of criticality are also observed by Haider *et al.* [8,9,10]. It is seen from Eqs. (13) and (18) that the amplitude of the solitary wave is three times larger than that of shock waves though the critical value for solitary and shock structures is the same. This figure also represents the above statement. Above the critical value the amplitude decrease with increasing the value of μ_n . It is also observed that increasing the value of μ_e the amplitude of the solitary and shock waves decreases and the critical value of μ_n shifted to higher value of it. Figure 3 shows the variation of the amplitude (ϕ_m) of solitary waves with σ_p and α , keeping the values $\mu_e = 0.05$, $\mu_n = 0.5$, and $u_0 = 0.1$. We have found a critical value of α which increases with increasing σ_p ; below (above) which the positive (negative) solitary and shock waves are found. We have seen that amplitude (ϕ_m) increases with increasing α below the critical value but decreases with α above the critical value. Figure 4 shows the variation of the width (δ_1) with μ_n and μ_e keeping the values $\sigma_p = 2$, $\alpha = 0.7$, and $u_0 = 0.1$; which indicates that δ_1 decreases with increasing μ_n and μ_e . Figure 5 shows the variation of the width (δ_1) with σ_p and α keeping the values $\mu_e = 0.05$ and $\mu_n = 0.5$, and $u_0 = 0.1$. We have found that δ_1 increases with increasing α and decreases with increasing σ_p .

6 Discussion

EPHIA solitary and shock structures have been analyzed in an unmagnetized plasma containing EPHI fluid with nonthermal distributed ENLIs and Maxwellian electron. We have investigated the basic features of amplitude, width, and temperature effects of electrons and ions. The results that we obtained from this investigation can be summarized as follows:

1. The amplitude of the solitary waves is three times higher than shock structures.
2. The amplitude of the solitary and shock waves for nonthermal ENLIs is proportional to the wave speed u_0 whereas the width is inversely proportional to it, which means that the profile of the faster soliton will be taller and narrower than slower one.
3. Depending on the parametric regimes both the solitary and shock waves might be associated with positive or negative potentials. There must be a critical value for nonthermal ENLIs and nonthermal parameter α keeping other parametric regimes constant, where the amplitude goes infinity.

4. Below (above) the critical value of ENLIs concentration the solitary and shock waves are associated with positive (negative) potentials. And also negative (positive) potentials are found below (above) the critical value of the nonthermal parameter (α).
5. Maxwellian electrons concentration makes the lower amplitude solitary and shock waves.
6. The width of solitary structures decreases with increasing α , σ_p , μ_n and μ_e ; i.e. the solitary waves will be spiky for the higher values of these.
7. The width of the shock waves is a linear function of the viscous coefficient, i.e. spiky profile of the shock waves are found for the lower value of the viscous coefficient.

We hope that our present investigations may be helpful to understand the effect of heavy ion on EPHIA solitary and shock waves in multi-component electron-ion plasmas in laboratory and astrophysical system like radio wave propagation in the ionosphere [12] and also give a guideline to further development in the relevant field.

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